

# Beam Physics Studies at Tokyo Tech

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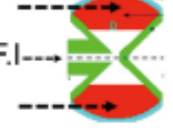
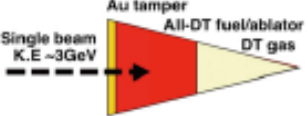

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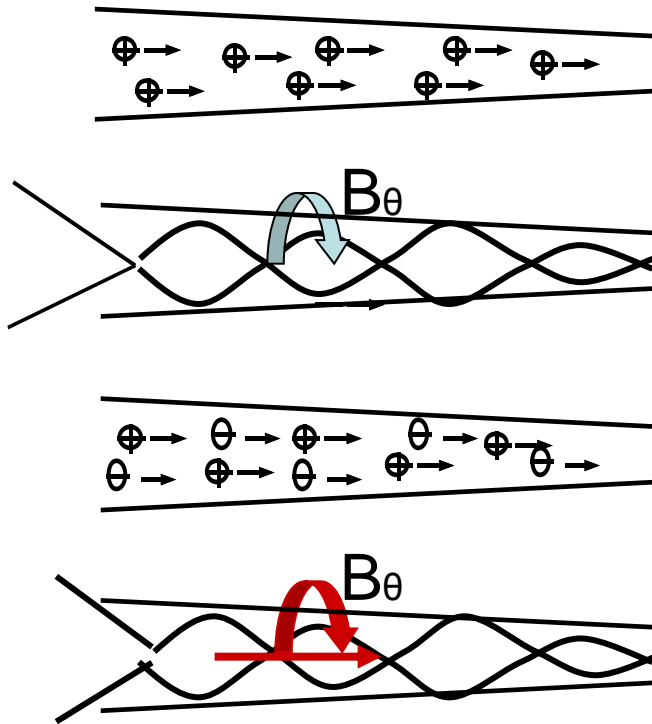
# HIF-targets cover a wide range of designs which request a variety of beam parameters

	Features	Issues
<b>Indirect drive –</b> <i>HS ign.</i> 	<ul style="list-style-type: none"> <li>• Integrated 2D designs exist</li> <li>• Ablation/burn physics on NIF</li> <li>• Natural two-sided geometry</li> </ul>	<ul style="list-style-type: none"> <li>• Lower drive efficiency</li> <li>• Lower gains, high driver energies</li> </ul>
<b>Direct drive X-target –</b> <i>Fast ign.</i> 	<ul style="list-style-type: none"> <li>• Inherent one-sided drive</li> <li>• High coupling efficiencies</li> <li>• Reduced stability issues</li> <li>• Potential for high yields (~GJ) and gains</li> </ul>	<ul style="list-style-type: none"> <li>• High gains require high densities under quasi-3D compression</li> <li>• Higher ion kinetic energies</li> <li>• High power, small focal spot beams needed for fast ignition</li> <li>• Driver concepts immature</li> </ul>
<b>Direct (+indirect) drive, tamped –</b> <i>Shock ign.</i> 	<ul style="list-style-type: none"> <li>• High coupling efficiencies (tamped ablation)</li> <li>• Simple targets</li> <li>• High gains consistent with low ion-kinetic-energies (~2-10GeV)</li> </ul>	<ul style="list-style-type: none"> <li>• Optimum ion species and energy</li> <li>• Two-sided (polar) geometry to be established**</li> <li>• High power beams needed for shock ignition</li> <li>• Stability to be confirmed</li> </ul>
<b>Direct drive, cylindrical compression –</b> <i>Fast ign.</i> 	<ul style="list-style-type: none"> <li>• Inherent one-sided drive</li> <li>• High coupling efficiencies</li> <li>• Simple targets</li> </ul>	<ul style="list-style-type: none"> <li>• Low gains, high driver energies</li> <li>• High ion kinetic energies</li> <li>• High power, small focal spot beams needed for fast ignition</li> <li>• Driver concepts immature</li> <li>• No U.S target design interest</li> </ul>

\*\*Will leverage present NIF PDD studies

7

# Beam transport scheme from final focus to the target is also an important issue



- Ballistic Transport
- Self-pinched Transport
- Ballistic with Electrons
- In pre-formed Plasma Channel

Possible chamber transport scheme depends on beam current and energy

# HIF-Accelerator system and target designs must be correlated closely

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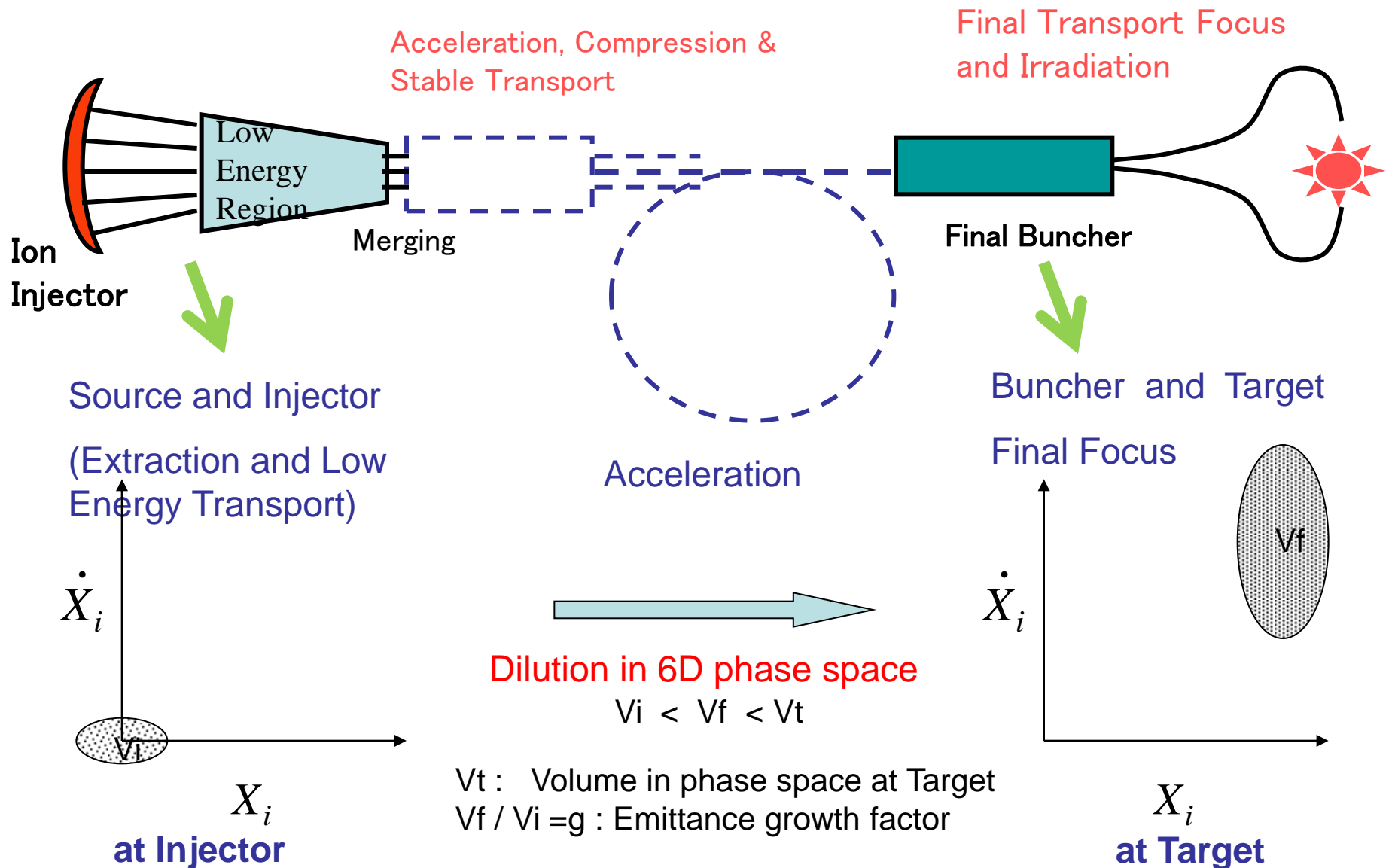
- Allows various target designs and chamber transport scheme
  - Direct or indirect
  - Cylindrical, spherical and X-targets
  - Chamber transport scheme
- Optimum ion species, its energy and beam number are still non-fixed
- We have to consider a rational accelerator scheme in the multidimensional parameter space

# We can standardize the problems with discussion of particles in the phase space

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- Necessary requirement that all of driver scheme must be met is;
  - The phase space density of particles must be less than the achievable phase density at the injector
- Goals of beam physics studies at Tokyo Tech are;
  - Quest for high-flux and low emittance injector
  - Investigate the phase density dilution (emittance growth) process from the injector to the target chamber
  - Discuss those beam dynamics in high power accelerators with laboratory-scale devices

# There are strong space charge issues in the ion injector and the final stage of high power ion accelerators



# We need a breakthrough for HIF injector

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High-flux ( $\sim 10^{-4} \text{ A} \cdot \text{sec}$  (  $10^{15}$  particles), and low emittance ( $V_i < V_t / g$  ) )

## Ion Extraction from Non-Stationary Moving Plasmas

### Expanding plasma

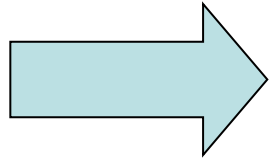
- Laser ablation plasma through strong axial field
- Plasma acceleration and gas-dynamic cooling
- Low temperature and high flux

Achievable beam flux and emittance (Phase space density)

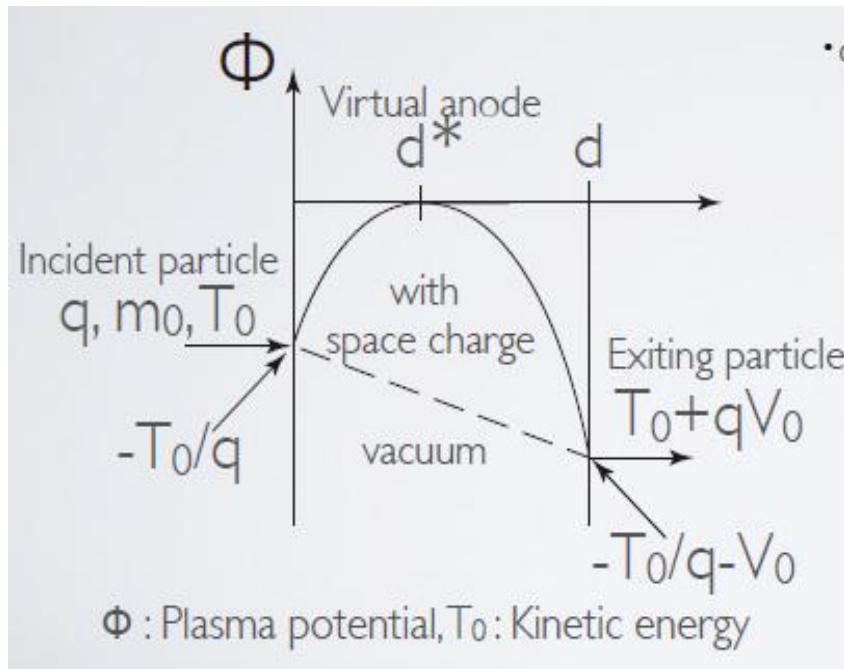
# Beam Dynamics during Ion Extraction from Moving Plasma

## Plasma Injection

$n_i, v_i, T_i, J_p$



## Acceleration Gap



Emission Surface  $d^*$

Potential Hump  $\phi_h$

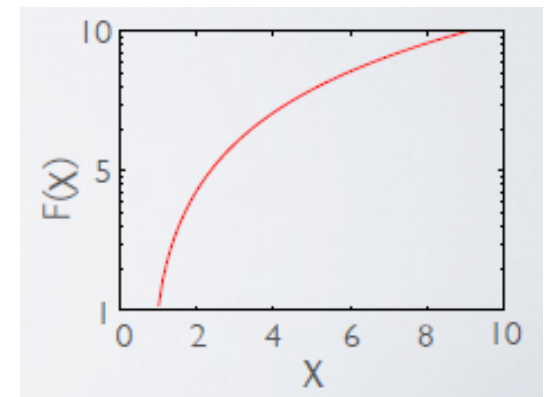
Enhancement  $F(\chi)$

## Beam Extraction

$J_i, \Delta E_i, \varepsilon_i$



$$J_i \propto J_{CL} \bullet F(\chi)$$

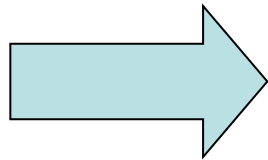


Enhancement Factor

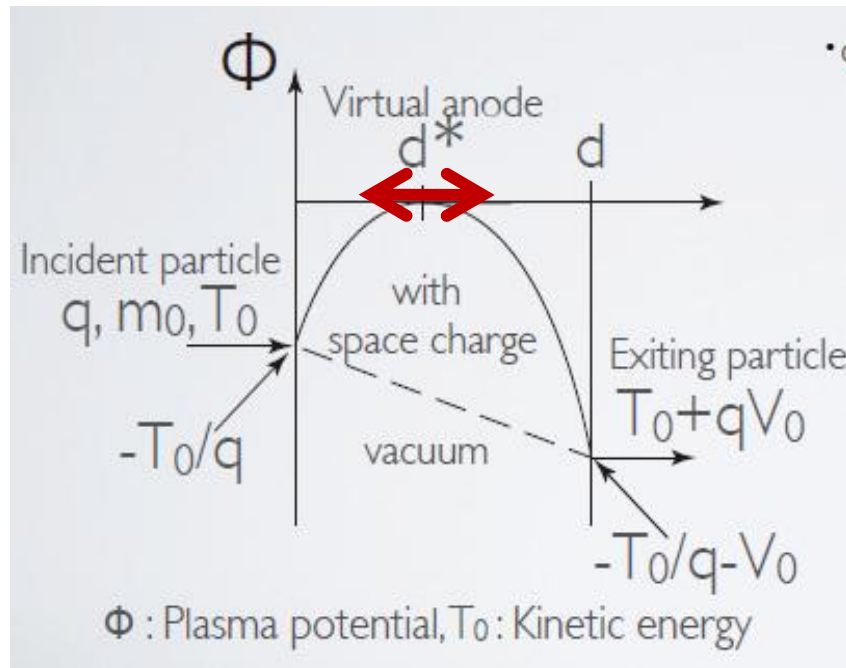
# Beam Dynamics during Ion Extraction from Moving Plasma

## Plasma Injection

$n_i, v_i, T_i, \underline{J_p}$



## Acceleration Gap



Emission Surface  $d^*$

Potential Hump  $\phi_h$

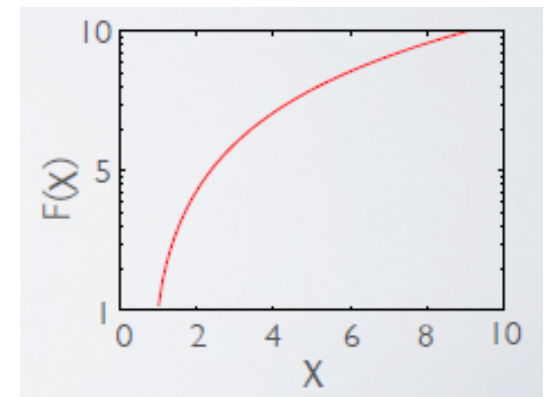
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## Beam Extraction

$\underline{J_i}, \Delta E_i, \epsilon_i$



$$J_i \propto J_{CL} \bullet F(\chi)$$

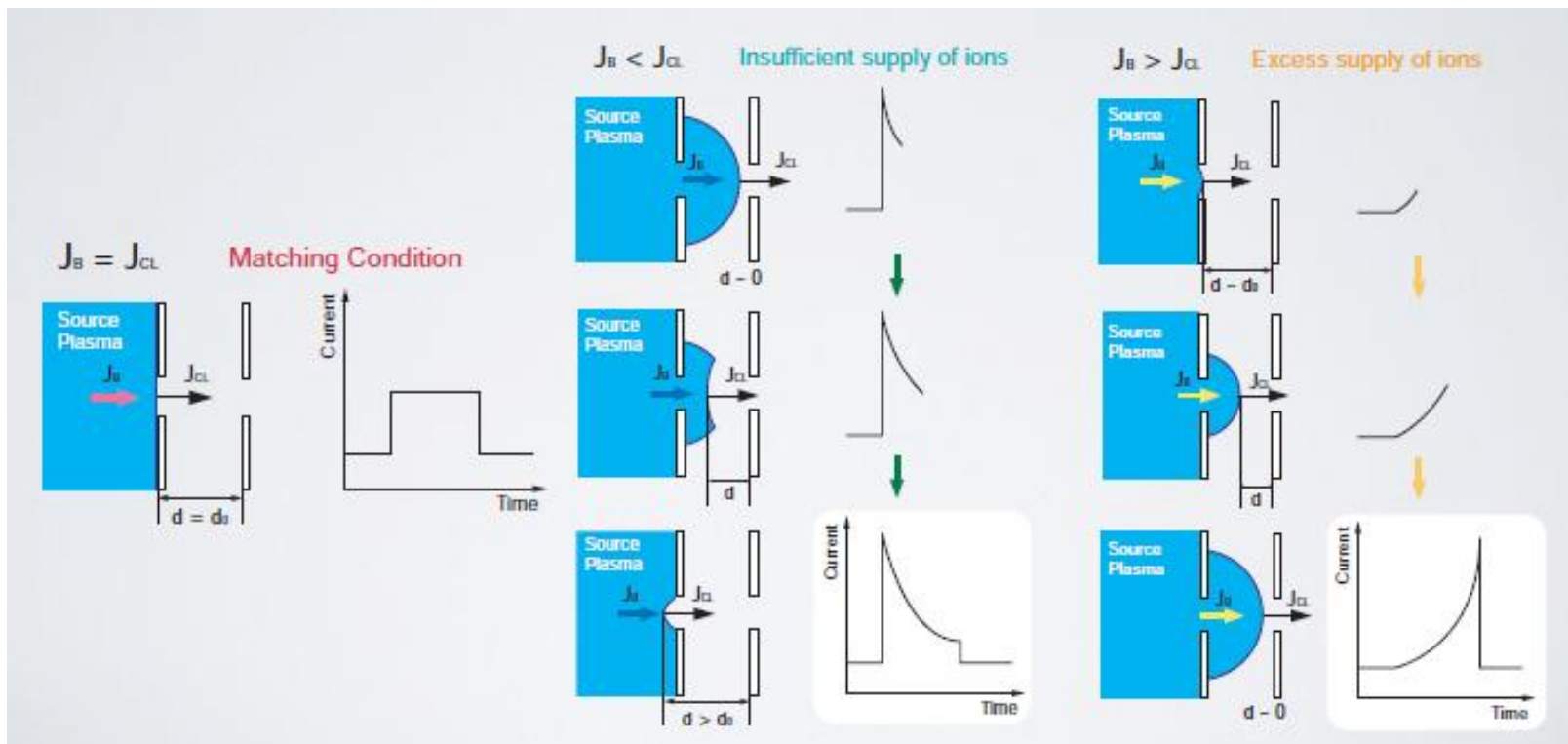


Enhancement Factor

# Example of Ion Extraction Process from a Moving Plasma

Moving and quasi-stationary source plasma  $\frac{mv^2}{2q} \ll V_0$

Behavior of the extraction gap is dominated by the parameter of source plasma



# High flux Ions are directly extracted from laser ablation plasma

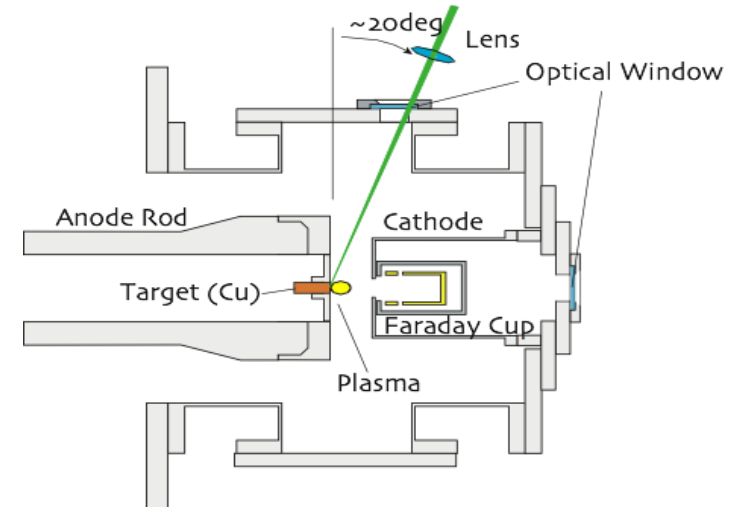
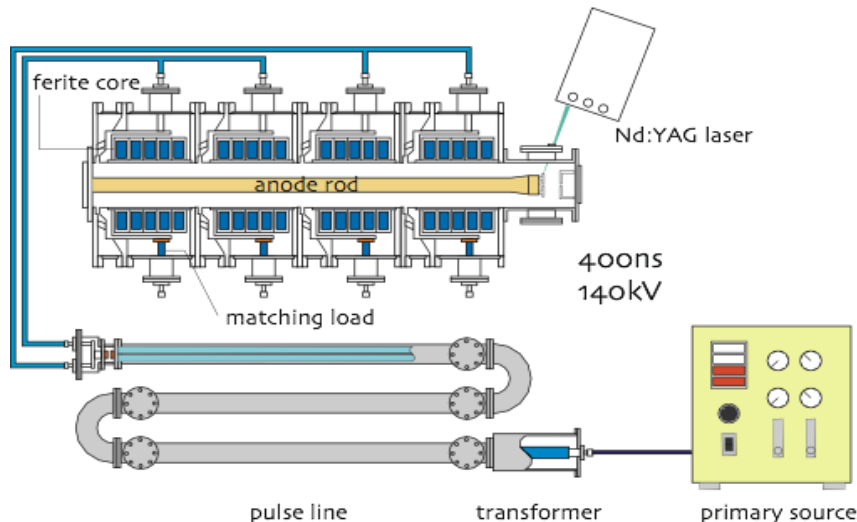
Direct ion extraction from drifting plasma

Cu Ions are extracted from 20mm gap

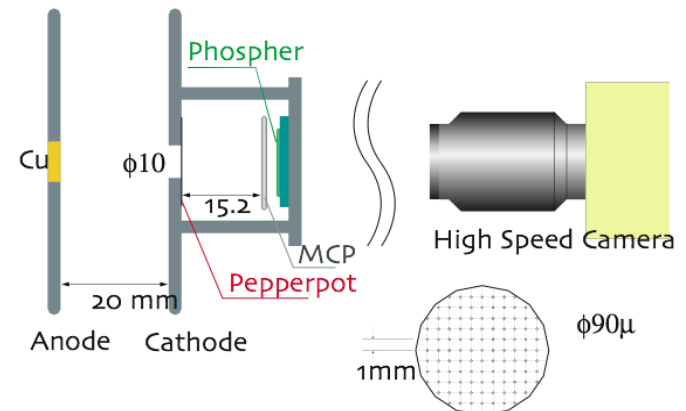
Beam current → Faraday Cup

Beam emittance → Pepper-pot image

A schematic view of laser ion source



Pepper-pot Emittance Measurement



# With properly adjusting the operating condition, high flux & low emittance ion beam can be extracted from ablation plasma

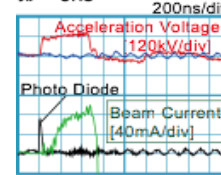
Can overcome the Bohm limit

Matching problem is overcome by controlling the ion supply close to the space charge limiting current of effective gap

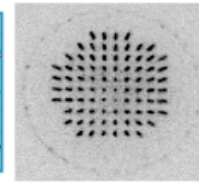
Beam bunch of Cu ions with 100mA/cm<sup>2</sup> level with emittance of  $0.25\pi$  mm·mrad and flat-top waveform was obtained

a) 33 mJ

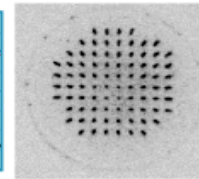
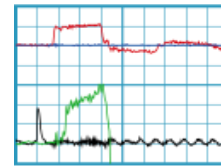
$t_d = 0$  ns



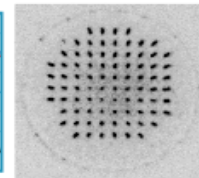
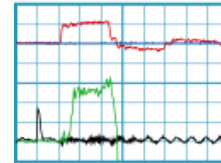
Beam Image



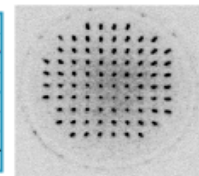
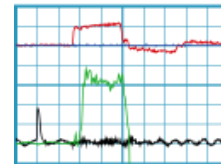
$t_d = 100$  ns



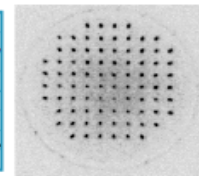
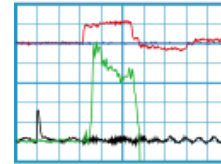
$t_d = 200$  ns



$t_d = 300$  ns

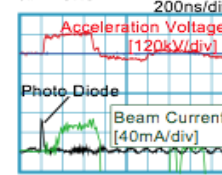


$t_d = 400$  ns

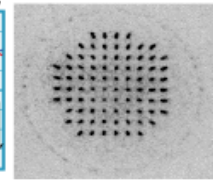


b) 28 mJ

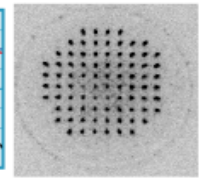
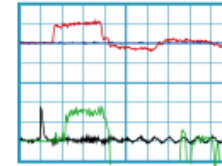
$t_d = 0$  ns



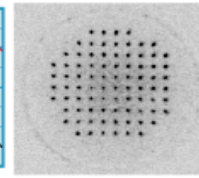
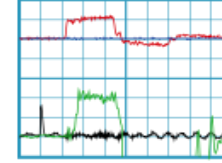
Beam Image



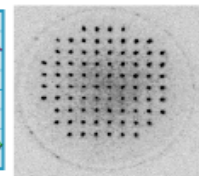
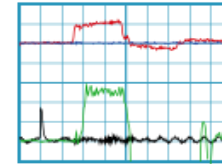
$t_d = 100$  ns



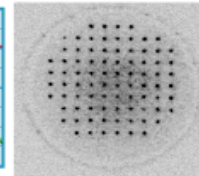
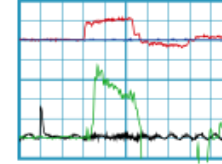
$t_d = 200$  ns



$t_d = 300$  ns



$t_d = 400$  ns



Matching condition

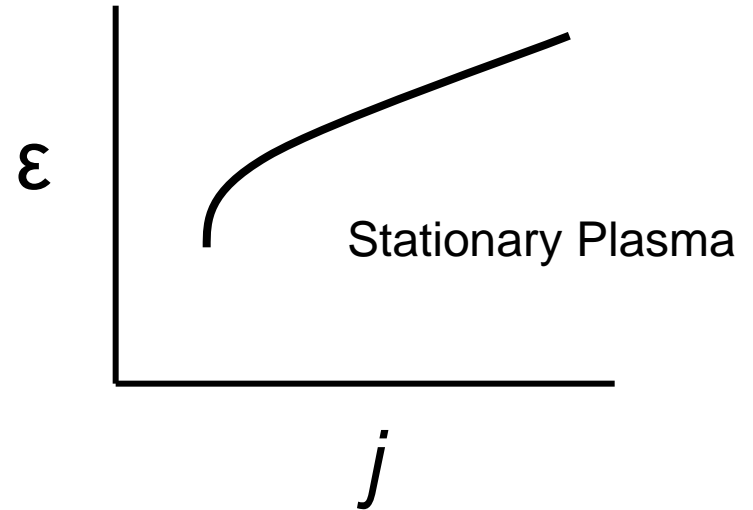


# Correlation between beam current and emittance for moving plasma source

- Flux from plasma source

$$j \propto Z e n_i u$$

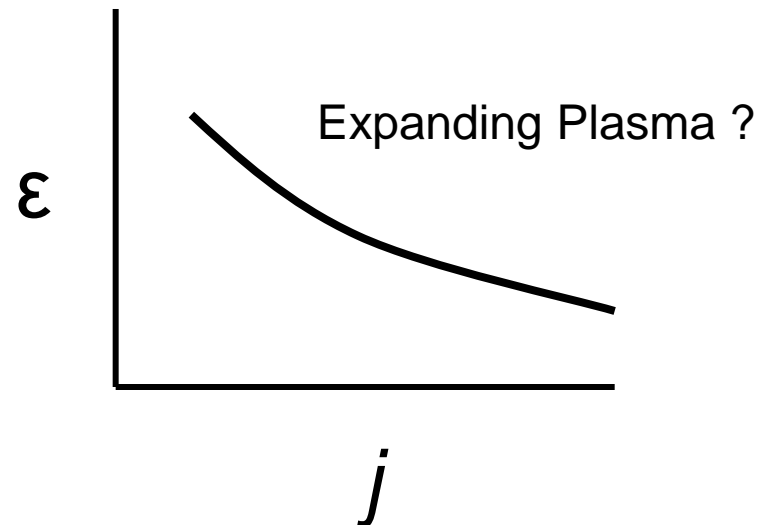
$$\varepsilon \propto T_i^{\frac{1}{2}}$$



- Stationary plasma

$$u \geq \left( \frac{k T_e}{m_i} \right)^{\frac{1}{2}}$$

$$j_B = e n_s v_0 = 0.6 e n_0 \sqrt{\frac{\kappa T_e}{m_i}} \propto \sqrt{T_e}$$



- Expanding plasma

$$u^2 = \frac{2\gamma R}{\gamma - 1} (T_0 - T_i)$$

$$j \propto \sqrt{(T_0 - T_i)}$$

# Moving (expanding) plasma source can increase the phase space density

Conventional injector (by J.Barnard)

$$\frac{dN}{dU_6} = \frac{N}{\Delta P_x \Delta P_y \Delta P_z \Delta x \Delta y \Delta z} = \frac{I}{q \gamma \beta^2 m^3 c^4 \epsilon_{nx} \epsilon_{ny} (\Delta P_z / P_z)} = \frac{\pi \epsilon_0}{2^{1/2} 9} \left( \frac{V_0}{m^3 q^3} \right)^{1/2} \frac{1}{k T_s (\Delta P_z / P_z) \alpha_B (1cm)^2}$$

Injector based on expanding plasma

$$\frac{dN}{dU_6} = \frac{I}{q \gamma \beta^2 m^3 c^4 \epsilon_{nx} \epsilon_{ny} (\Delta P_z / P_z)} = \frac{q n v_d A}{q \frac{2qV_0}{m} m^3 \left( \frac{kT_s}{m} \right) (\Delta P_z / P_z)} = \frac{A}{2qkmV_0} \frac{n v_d}{(\Delta P_z / P_z) T_s} \propto \frac{(T_0 - T)^{1/2}}{(\Delta P_z / P_z) T^{\frac{\gamma-2}{\gamma-1}}}$$

$$n \propto c T^{\frac{1}{\gamma-1}}, \quad v_d \propto \left( \frac{\gamma}{\gamma-1} R(T_0 - T) \right)$$

Expansion (cooling & acceleration) ➡ High-flux and high phase space density

# Summary of the injector study

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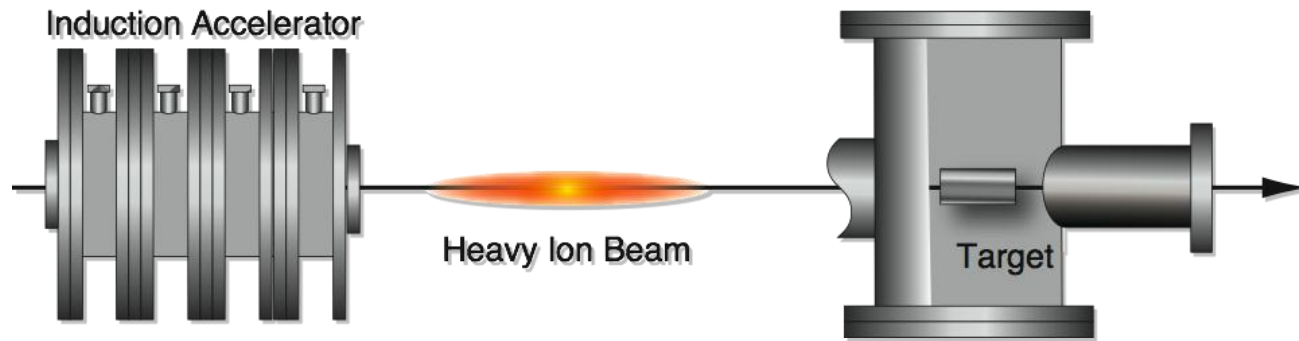
## Summary

- We discussed ;
  - Ion extraction process from laser ablation (moving) plasma
  - Behavior of laser ablation plasma through axial magnetic field

## Future Plans

- Modify the drifted-Maxwellian particle distribution
- Discuss the achievable flux, phase space density and pulse length of the ion injectors

# Issues for Longitudinal Bunch Compression Experiments



- Bunch Compression Experiments using Induction Voltage Modulator
- Emittance Growth accompanied by the Beam Manipulation
- Coupling Effects between Longitudinal and Transverse Modes
- Collective Effects during the Bunch Compression

## Influence Factors of Compression Ratio (Emittance Growth)

- Accuracy of Modulation Voltage
- Longitudinal Emittance at Injector
- Space Charge Effect
- Collective Effect

# Beam Dynamics during Bunch Compression in the Final Stage

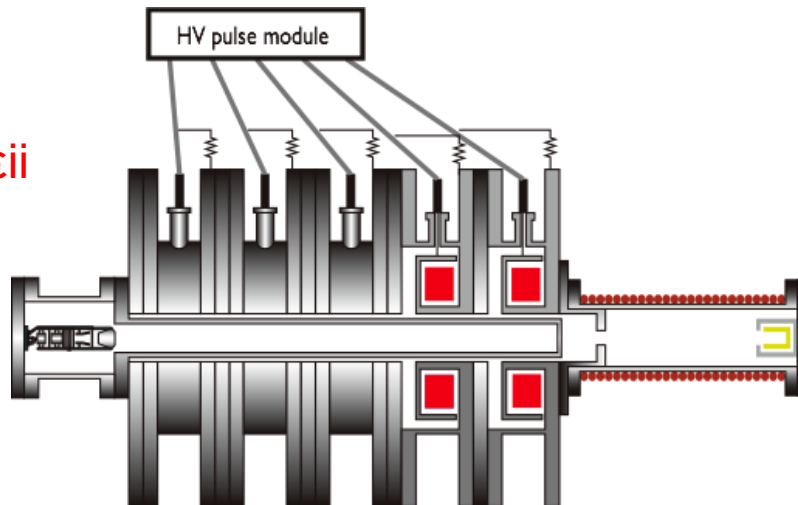
**Beam Injection**

**Beam Bunching and Transport**

**Beam Extraction**

$E_i, \Delta E_{if}, J_{ii}, T_{ii}, \epsilon_{ii}$

$E_{if}, \Delta E_{if}, T_{if}, J_{if}, \epsilon_{if}$



Space Charge Field

Collective Effect

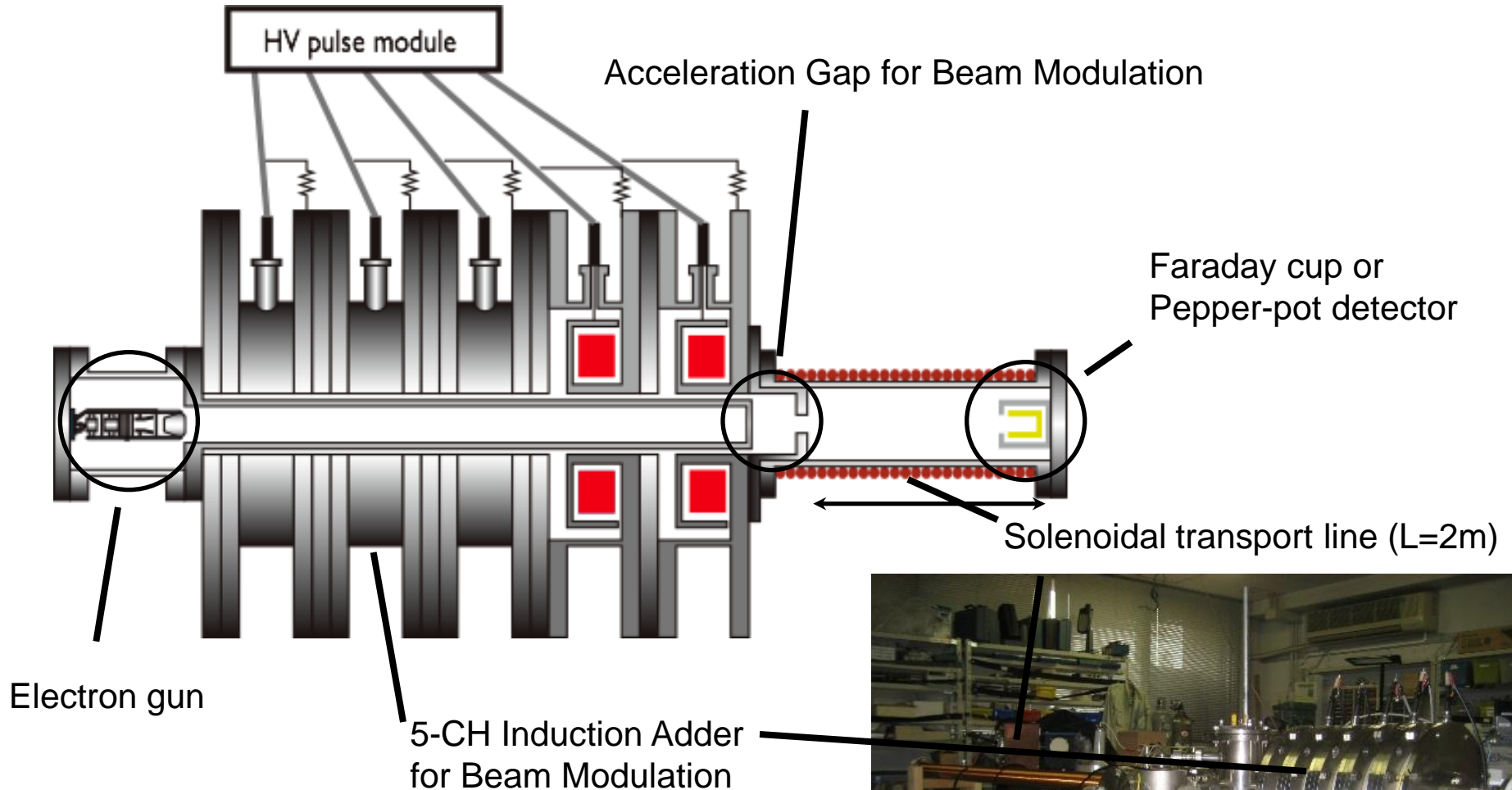
Longitudinal-Transverse Coupling

$$F(\chi) = \frac{J_{if}}{J_{ii}} \propto f(\epsilon_{if})$$

**Emittance Growth**

**Compression Ratio**

# Arrangement for Bunch Compression Experiments

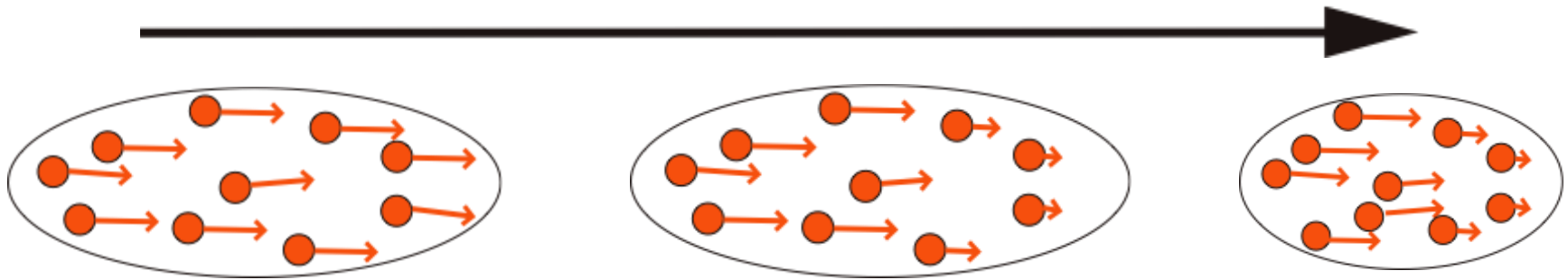


## Strategy

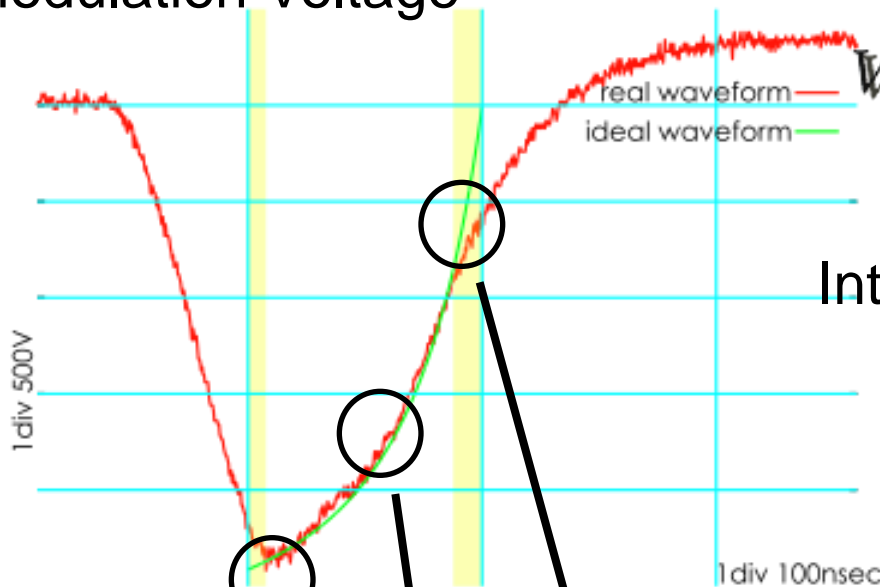
- FET-based voltage driver for Induction Adder
- E-beam using Grid-controlled e-gun
- Solenoidal transport line



# Analytically Derived Manipulation-Voltage for Longitudinal Compression



Modulation Voltage



Intended Voltage at Beam buncher

$V_0$  : Initial Beam-Voltage

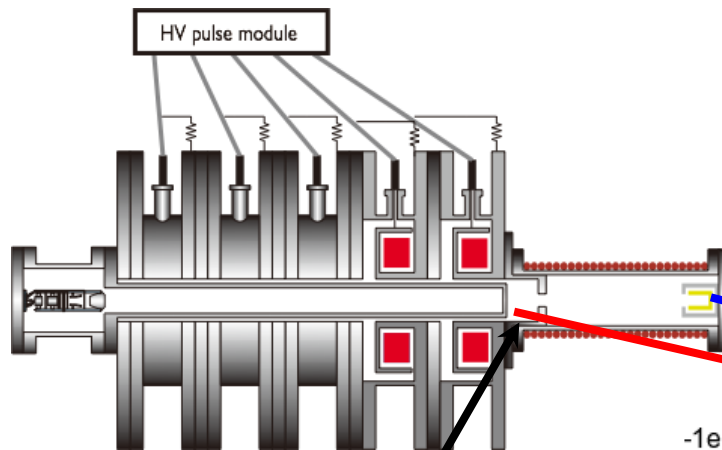
$T$  : Initial Pulse Width

$L$  : Transport Distance

$$V_{dec}(t) \equiv \frac{m_e}{2q_e} \cdot \frac{1}{\left( \sqrt{\frac{m_e}{2q_e V_0}} + \frac{T-t}{L} \right)^2} = V_0$$

Beam front  
Beam middle  
Beam end

# Typical Result of PIC Simulation of the e-Beam Transport

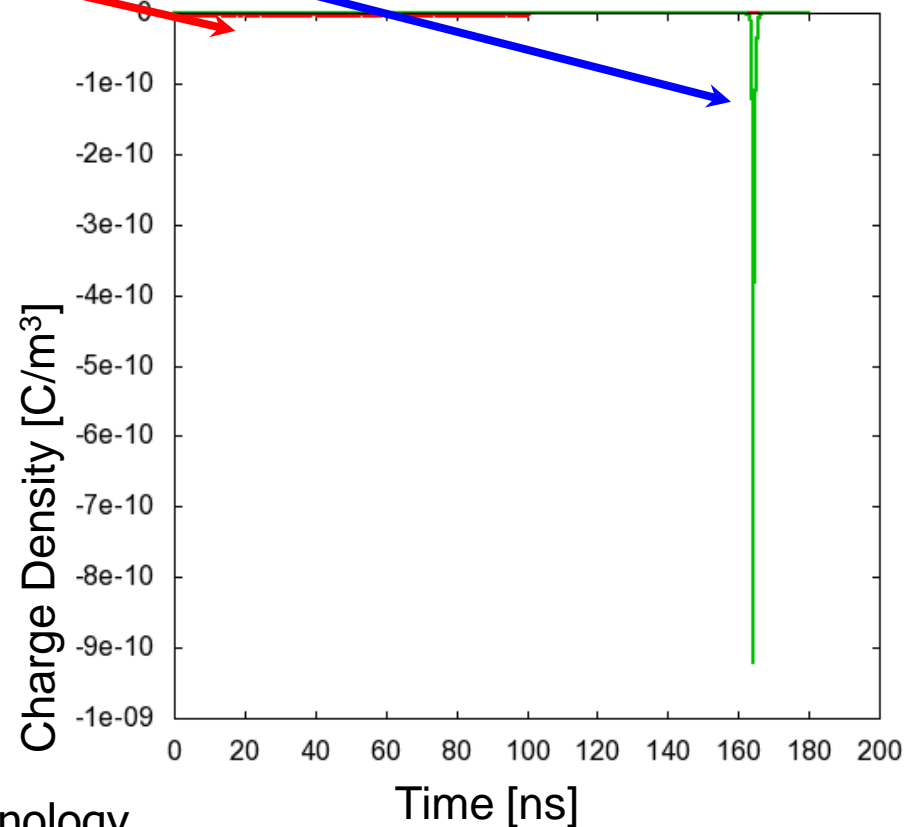


1D PIC Calculation

Applied Voltage

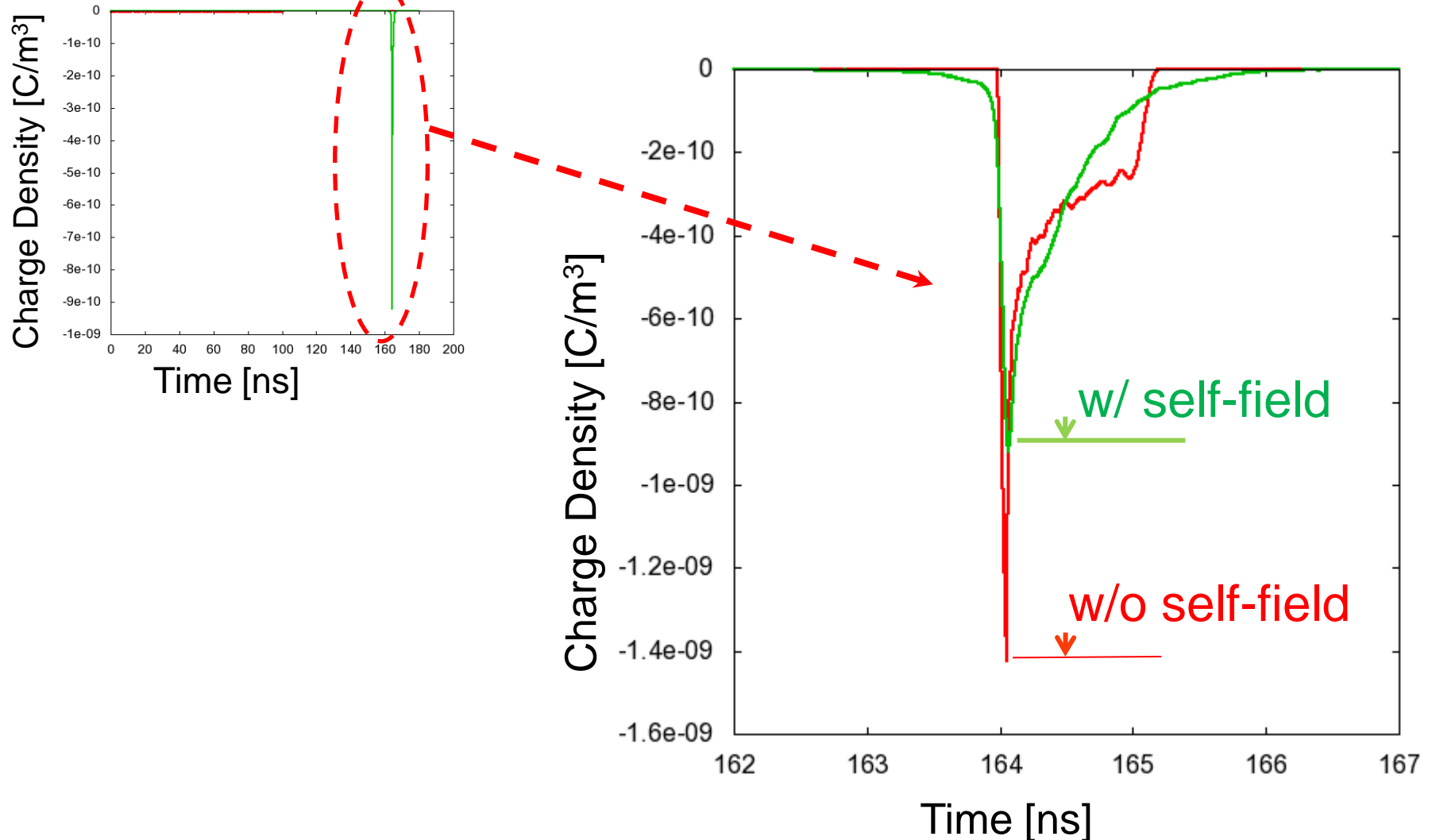
$$V_{dec}(t) = \frac{m_e}{2q_e} \cdot \frac{1}{\left( \sqrt{\frac{m_e}{2q_e V_0}} + \frac{T-t}{L} \right)^2} - V_0$$

$V_0 = 2.8 \text{ kV}$  @  $L = 2 \text{ m}$  during  $T = 100 \text{ ns}$

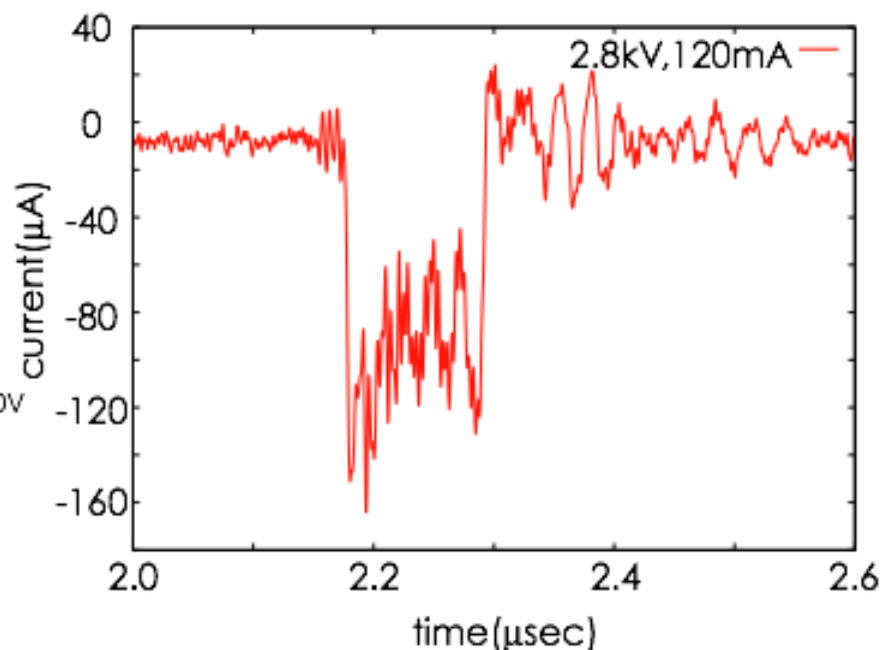
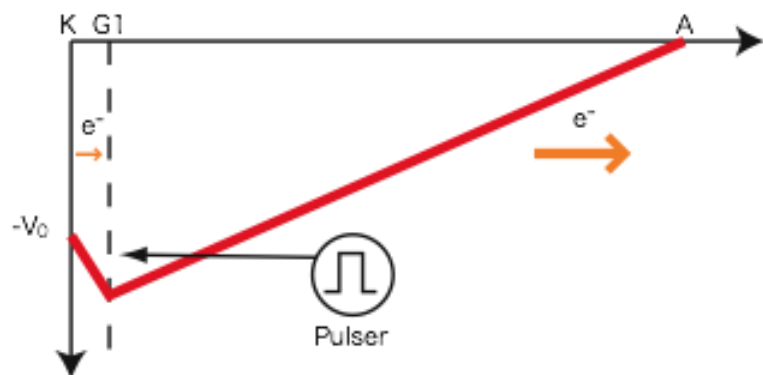
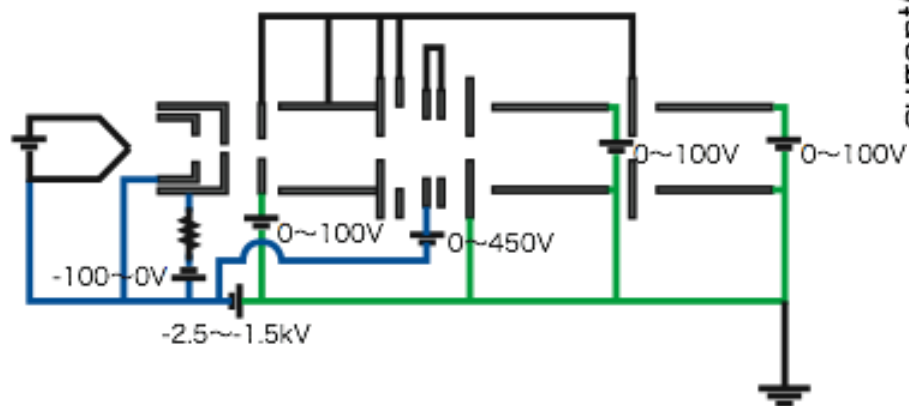
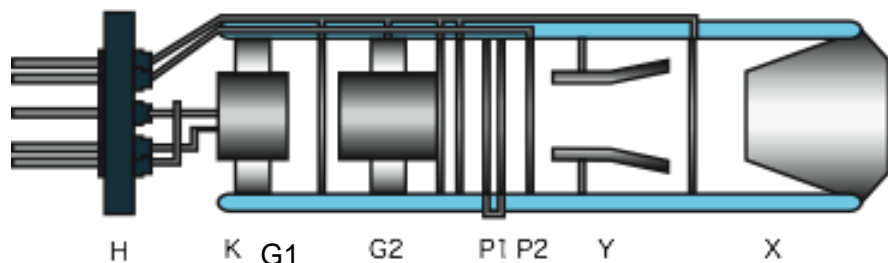


# PIC Simulation indicates an Effect of Space Charge Field on the Bunch Compression in Our Experimental Condition

## Evolution of Charge Density of e-Beam Transport with Compression



# Grid-controlled Electron Gun for Low Temperature Beam Formation

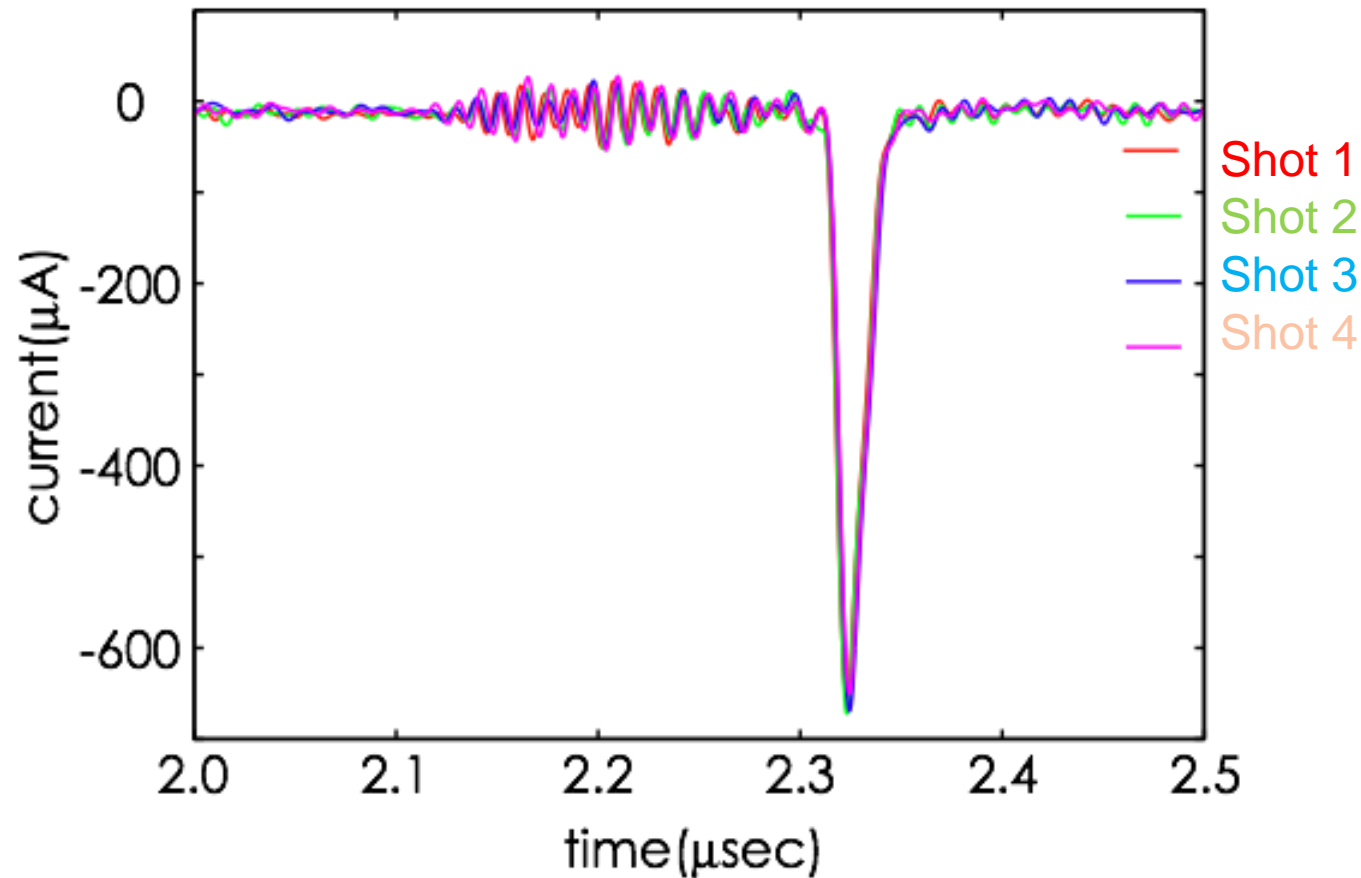


Typical Waveform of e-Injector Current

## Electron Gun

- Thermo-electric Emission
- Statically Biased Electrodes
- Grid-controlled

# Reproducibility of the Compressed Waveform

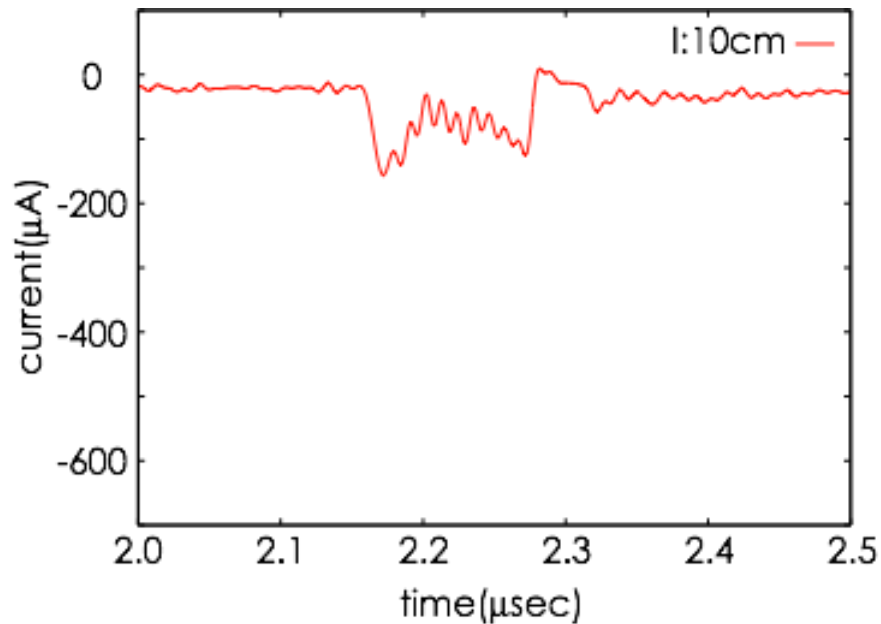


Successive 4-shots traces were overlapped

Reproducibility was vastly improved by the Grid-controlled electron gun and FET voltage modulator

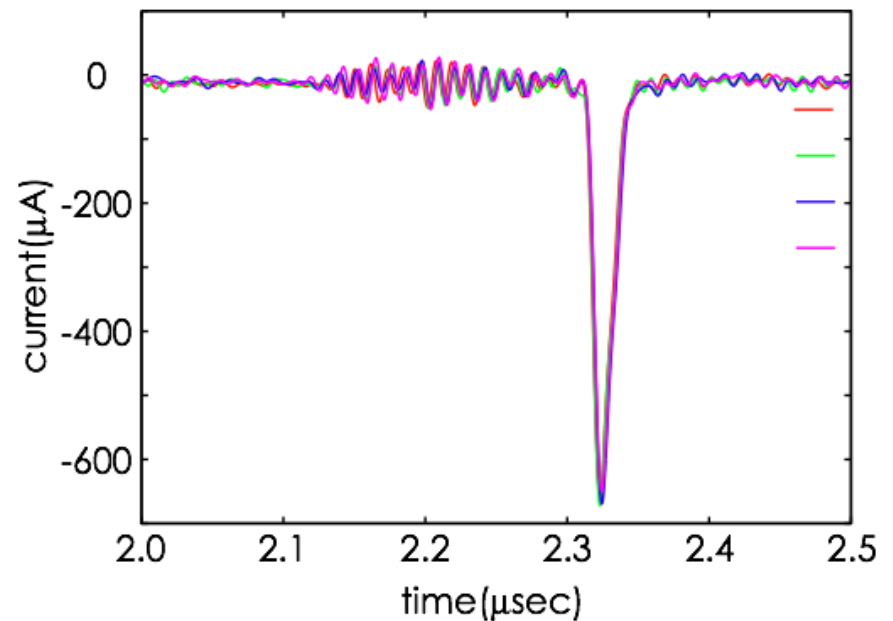
# Waveforms of Beam Current at Injection and at the Destination

Injection 2.8keV, Pulsed Beam  $T=100\text{ns}$ ,  $B \simeq 0.03T$



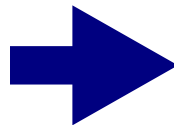
Current at beam injector

$$\begin{aligned} I_{peak} &\approx 100\mu A \\ \Delta t &\approx 100nsec \\ J &\approx 3.2mA/cm^2 \end{aligned}$$



Beam current at destination (L=2m)  
Successive 4-shots were over lapped

$$\begin{aligned} I_{peak} &\approx 670\mu A \\ \Delta t &\approx 15nsec \\ J &\approx 21mA/cm^2 \end{aligned}$$



# Summary and Future Plans of the Beam Compression Experiments

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## Summary

- We developed ;
  - FET-driven voltage adder for longitudinal bunch compression
  - Grid-controlled quasi-static electron gun
  - 2m solenoidal transport line
- Beam bunch was compressed with factor 6-7 and the process has good reproducibility

## Future Plans

- Faster FET switch for more precise modulation waveform
- Stronger magnetic field for suppression of transverse beam motion
- Discuss transverse-longitudinal coupling and/or collective effects on the emittance growth

# Concluding Remarks

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- HIF and/or HED drivers need beam manipulation in strongly space charge dominated region, in particular, in the ion injector and the final stage of the accelerator
- For high-flux and low emittance ion extraction:
  - We are investigating ion extraction processes from moving plasma
- For study on the bunching-beam dynamics:
  - We have started a scaled experiment using a controllable/reproducible device composed of grid-controlled e-beam source and FET-driven induction buncher
- Goal of our beam physics studies is to know the margin of emittance growth from the injector to the target